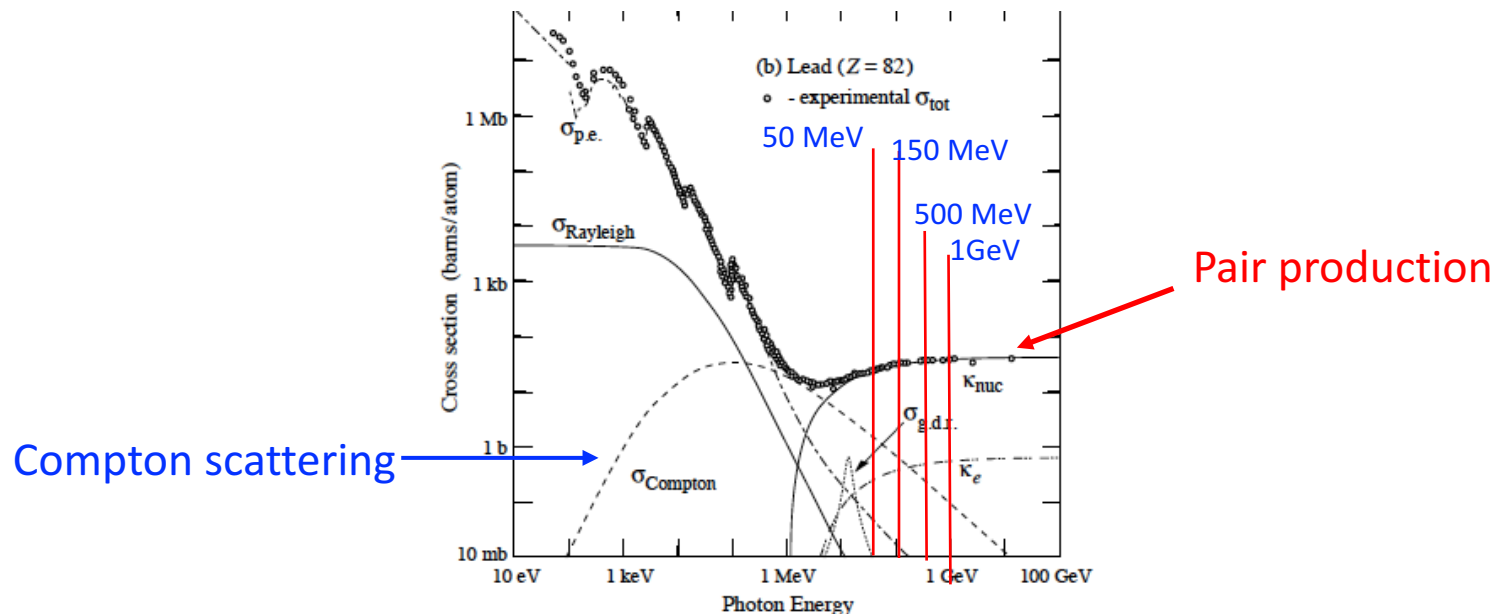


# Update on photon analysis

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2021.02.04

# Photon cross section



**Figure 27.15:** Photon total cross sections as a function of energy in carbon and lead, showing the contributions of different processes [48]:

- $\sigma_{\text{p.e.}}$  = Atomic photoelectric effect (electron ejection, photon absorption)
- $\sigma_{\text{Rayleigh}}$  = Rayleigh (coherent) scattering—atom neither ionized nor excited
- $\sigma_{\text{Compton}}$  = Incoherent scattering (Compton scattering off an electron)
- $\kappa_{\text{nuc}}$  = Pair production, nuclear field
- $\kappa_e$  = Pair production, electron field
- $\sigma_{\text{g.d.r.}}$  = Photonuclear interactions, most notably the Giant Dipole Resonance [49]. In these interactions, the target nucleus is broken up.

At low energy it is seen that the photoelectric effect dominates.

At 50-150 MeV, the pair production dominates, although Compton scattering, photoelectric effect contribute.

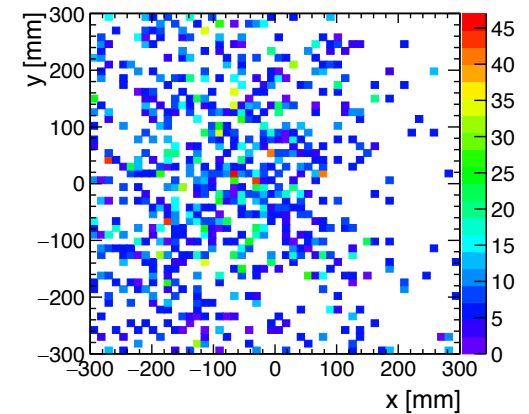
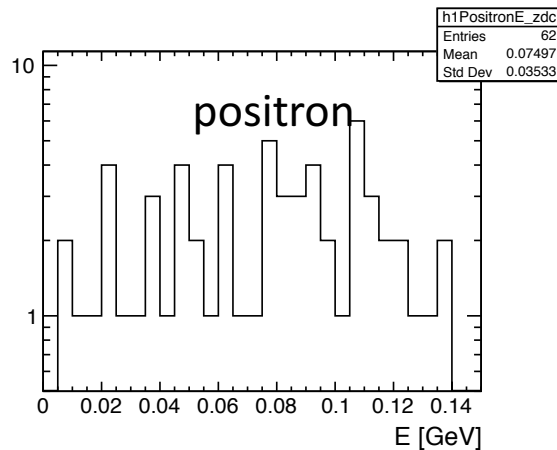
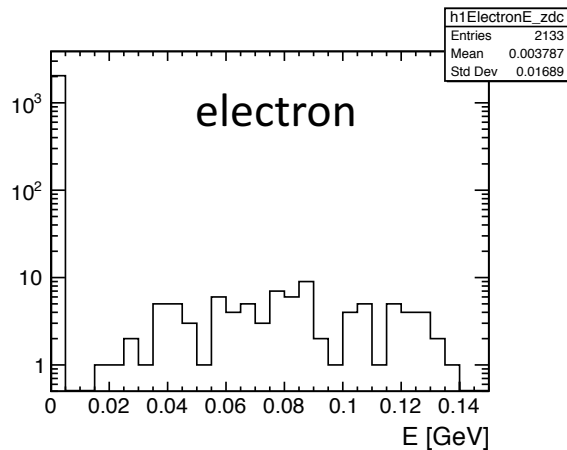
At 150-500 MeV and 500 MeV to 1 GeV, the increasing domination of pair production as the energy increasing, the Compton scattering contribution decreasing.

# Photon gun (0-5mrad)

Photon gun: theta 0-5mrad

10k events

	50-150MeV	150-500MeV	0.5-1GeV
Photon Convert	3232	3354	3474
Convert rate	32.32%	33.54%	34.74%



# BeAGLE run

1200 events

$e + \text{Pb} \rightarrow e' + \text{J}/\psi + \text{X}(\text{p}, \text{n}, \gamma)$

18x110 GeV

## Cut1:

- no neutron in ZDC

## Cut2 :

- Cut1 + no photon  $E > 50 \text{ MeV}$  in ZDC

## Cut3:

- Cut2 + no proton in Roman Pots

## Cut4:

- Cut3 + no proton in off-energy detector

## Cut5:

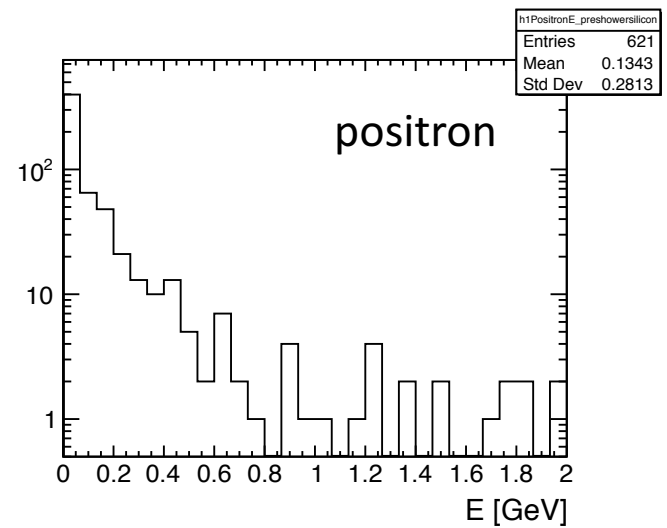
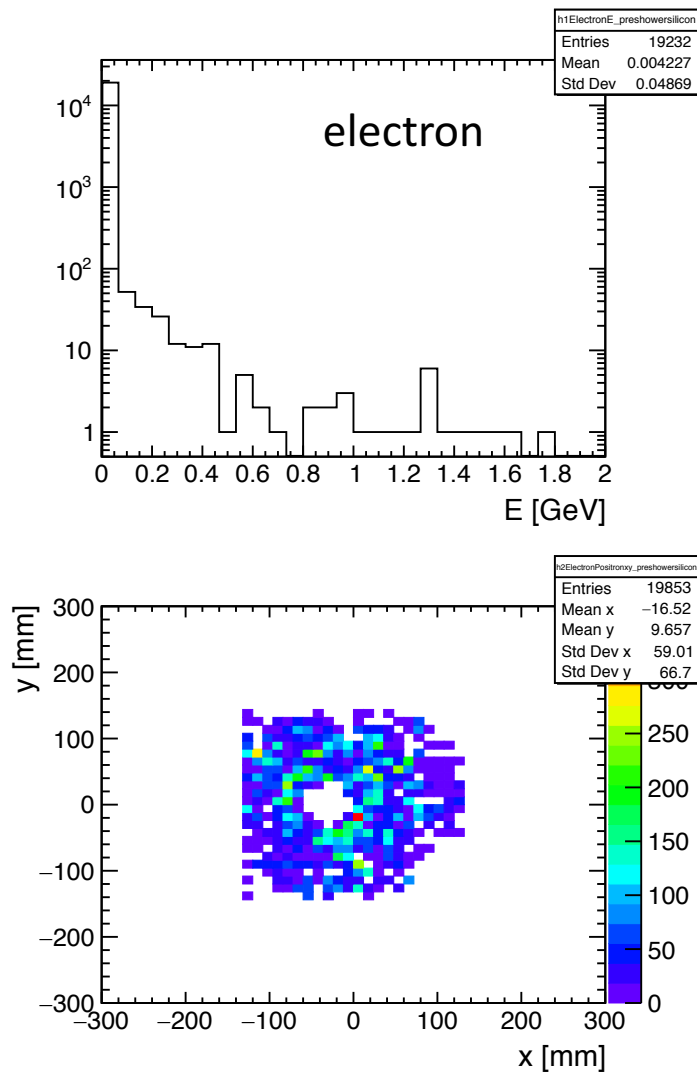
- Cut4 + no proton in B0

## Cut6:

- Cut5 + no hits in preshower

Survived event count	
Total events	1200
Cut1	205
Cut2	143
Cut3	143
Cut4	141
Cut5	139
Cut6	60

# BeAGLE run

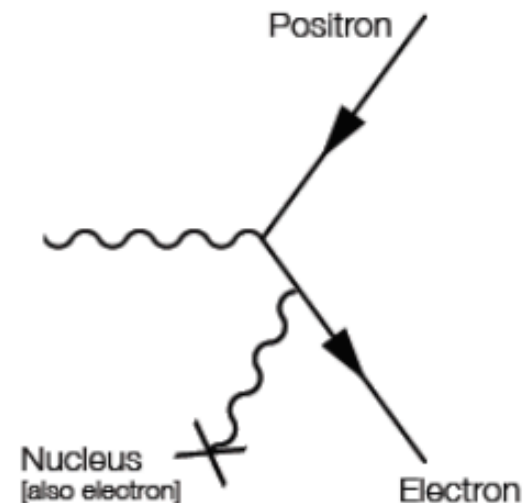
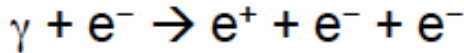
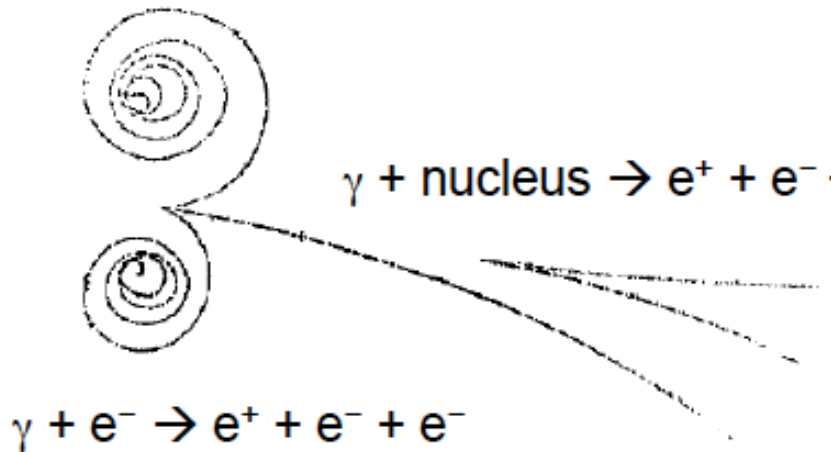
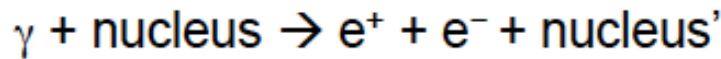
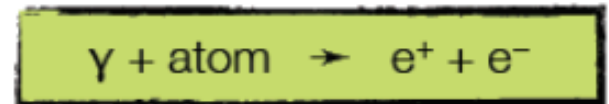
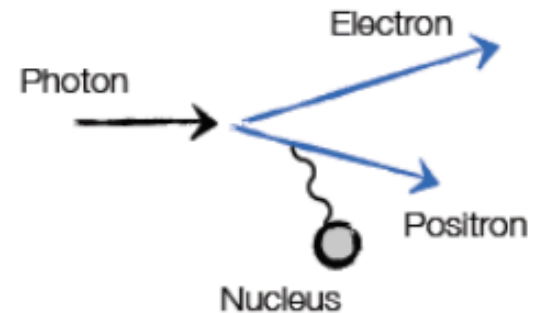


Back up

# Pair production

Minimum energy required for this process  
 $2 m_e c^2 + \text{Energy transferred to the nucleus}$

$$E_\gamma \geq 2m_e c^2 + \frac{2m_e c^2}{m_{\text{Nucleus}}}$$



# Compton scattering

Best known electromagnetic process  
(Klein–Nishina formula)

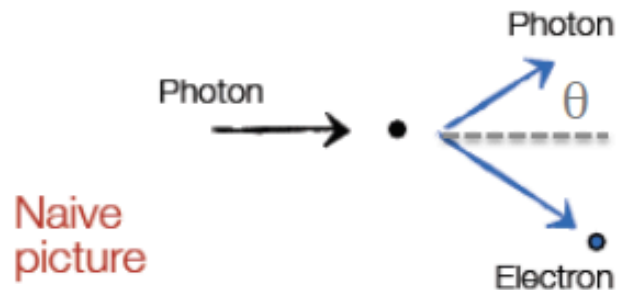
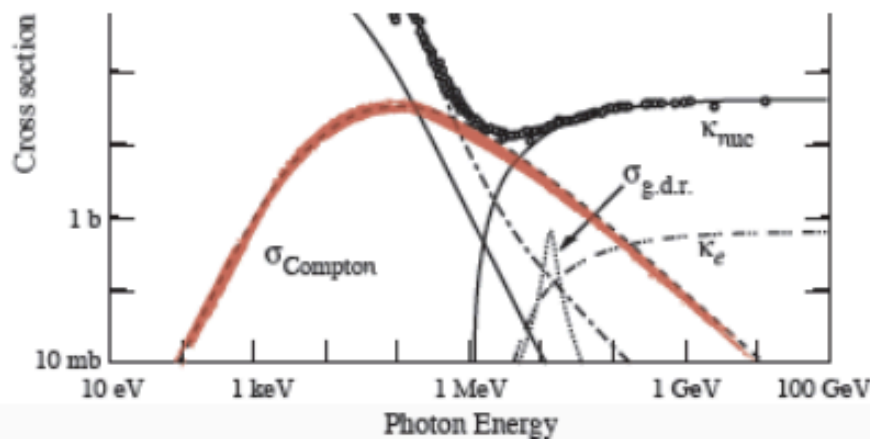
for  $E_\lambda \ll m_e c^2$   $\sigma_c \propto \sigma_{Th}(1 - 2\varepsilon)$

Thompson cross-section:

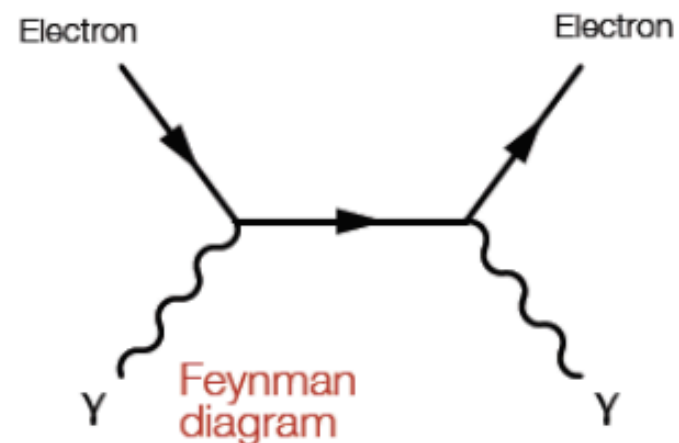
$$\sigma_{Th} = 8\pi/3 r_e^2 = 0.66 \text{ barn}$$

$$\varepsilon = \frac{E_\lambda}{m_e c^2}$$

for  $E_\lambda \gg m_e c^2$   $\sigma_c \propto \frac{\ln \varepsilon}{\varepsilon} Z$



$$\gamma + e^- \rightarrow (\gamma)' + (e^-)'$$





# Photoelectric effect

From energy conservation:

$$E_e = E_\gamma - E_N = h\nu - I_b$$

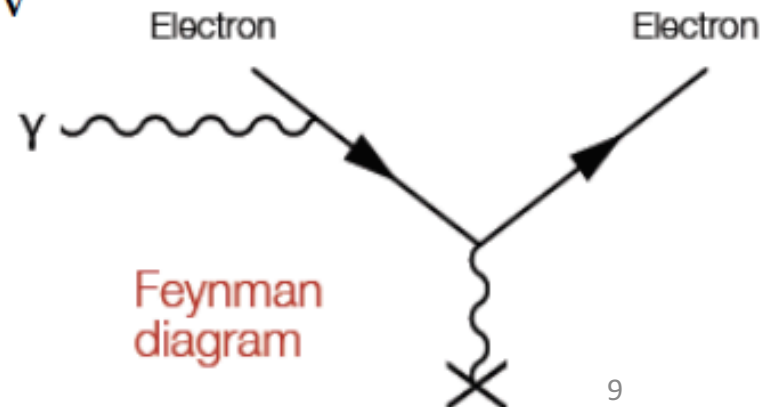
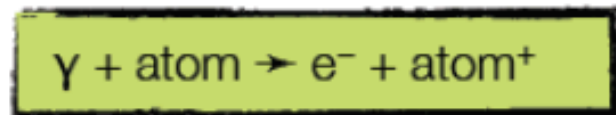
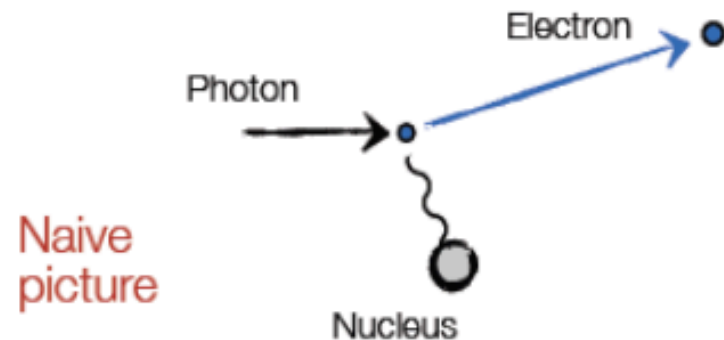
$I_b$  = Nucleus binding energy  
introduces strong  $Z$  dependence

Cross-section largest for  $E_\gamma \approx$  K-shell energy  
Strongest  $E$  dependence for  $I_0 < E_\gamma < m_e c^2$

$$\sigma_{ph} = \alpha \pi a_B^2 Z^5 (I_0 / E_\gamma)^{7/2} \quad \begin{matrix} a_B = 0.53 \text{ \AA} \\ I_0 = 13.6 \text{ eV} \end{matrix}$$

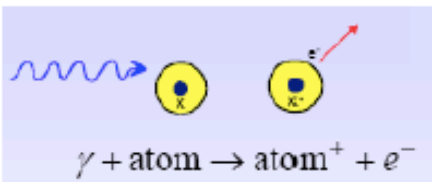
$E$ -dependence softer for  $E_\gamma > m_e c^2$

$$\sigma_{ph} = 2\pi r_e^2 \alpha^4 Z^5 (mc)^2 / E_\gamma$$

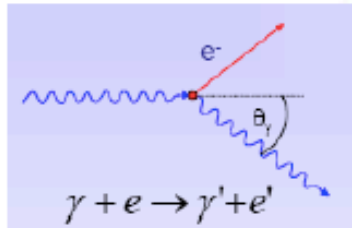


# Electromagnetic interactions

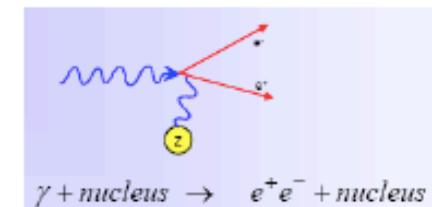
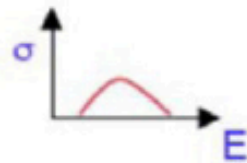
Gammas



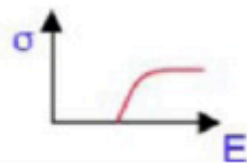
- Photoelectric effect



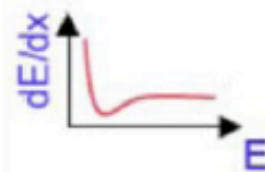
- Compton effect



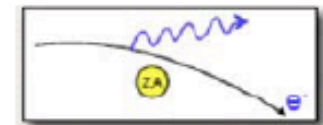
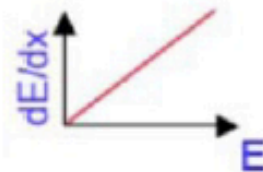
- Pair production



- Ionisation



- Bremsstrahlung



Electrons